24. SEQUENCE EFFECTS IN FATIGUE CRACK GROWTH DESCRIBED ON A CONTINUUM MECHANICS BASIS

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When regarding constant amplitude (C.A.) loading as the simplest type of cyclic loading (Fig. 1) common knowledge is that crack growth is a continuous process with an unequivocal law – for fixed environmental conditions depending mainly upon the effective stress intensity range, $\Delta K_{\rm eff}$, due to the load range above the crack opening level. The hope was that any kind of irregular loading could be interpreted as a series of C.A. cycles of varying magnitude. Crack length would then be the sum of these C.A. crack growth increments. Unfortunately, that is not the case. Cracks seem to grow as they want to do.

In the authors' opinion the crack growth law is valid furtheron. However, the real ΔK_{eff} development differs from that which would have been expected from considering the single cycles. ΔK_{eff} is governed by a kind of structural memory of the preceding load cycles. With other words: The rate of cyclic crack growth depends upon load sequence effects / 1 / which are due to the elasto-plastic response of materials / 2 /.

Two of the most typical load sequence effects are schematically depicted in Fig. 2. First, due to the jump from low to high loading this sequence shows accelerated crack growth. "Accelerated" means that the actual crack growth is faster than that of pure high load cycling without the preceding low cycles. Second, the load returning to the previous load level causes retarded crack growth – what means that the actual crack growth can be much slower than the crack growth under pure low load cycling. How much depends upon the height of the load jump and some other parameters. The schematic drawing gives the additional information that the life increasing retardation effect generally is one order of magnitude greater than the corresponding acceleration effect.

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Each irregular load-time function should be resolvable into discrete occurences of a few basic load sequences with known crack growth behaviour so that we are able to make realistic life predictions of structural components under service conditions.

Following this concept, in the past the authors investigated crack growth by means of the Fatigue Fracture Mechanics theory /3, 4 / and reported on computations of the elasto-plastic stress-strain field at a crack tip from load reversal to load reversal. The case studies brought insight into the continuumechanics background of the fatigue crack growth process which is mainly governed by crack closure. The need at the present is to find out what are the independent variables of the load sequence effects and what are the pure material parameters. In the present paper a rough approach to only one of several load sequence problems is tried. In order to make the results translucable it is based on some engineering assumptions rather than on more accurate crack closure solutions.

We restrict ourselves to the special case named "Single Peak Overloading" including a possible subsequent peak underload. Its typical crack growth curve is depicted on the right hand side of Fig. 3, the first derivation, i.e. the crack growth rate, on the left hand side. Where the sequence or load interaction effect ceases—that is where the actual crack growth rate reaches C.A. values—there is the point at which the delay or retardation effect can be quantified. The total number of affected cycles N_{CP} is split up into the number of cycles N_{CA} which would have been obtained without the peak load, plus the so-called number of delay cycles $N_{\text{D}} \neq 5$. N_{SP} and N_{D} shall be the descriptive quantities for the single peak overload response of the material. To formulate their dependence on the load sequence and the material parameters is the intention of the paper.

With Fig. 4 the field of phenomenogical considerations is left. You find compared the Fatigue Fracture Mechanics model / 4 / on the left hand side and on the right hand side the simplified model used here for establishing an approximate closed solution. The da/dN plot was rotated to the more common kind of representation with the crack length on the abscissa.

The refined theory / 4 /models the real behavior in thin sheets fairly well. Individual examples of the problem group concerned can be solved with this tool. However, it must be said that this way is too complicated for general use. In the figure you find the size-development of the plastic zone with crack length increase and the $\Delta K_{\mbox{eff}}$ -development which corresponds to the da/dN development the typical behaviour of which is in accordance with the experimental findings. Though, actually, theory / 2 / and most experiments (e.g. / 6 /) indicate that the interaction is not yet finished when the plastic zone strikes the peak-load-induced elasto-plastic boundary, the simplified model works with this endpoint. This deviation can be proved to have a minor influence on the model accuracy. A further simplication concerns the $\Delta K_{\mbox{eff}}$ -function which we assume to consist of two straight lines to and from the minimum. This was done to achieve a simple closed integration.

The relationship for N_{SP} shown in Fig. 5 is derived by means of the crack growth law $^{1)}$ and the simplified $\triangle K_{eff}$ function. When analyzing the integral solution we observe that the second term guided by $\triangle a_M$ can be neglected since, compared to the first term, it inherits the product of two small quantities.

As a result the total number of cycles is a function of the exponent of the C.A. crack growth law, a function of $\triangle a_D$ which can be easily calculated, and of \bigvee_D which is the ratio of two $\triangle K_{eff}$ values. Equivalently, \bigvee_D can be written as the ratio of the stabilized crack growth rate to the minimum crack growth rate. The problem, however, is not yet solved as the minimum growth rate value still is an unknown quantity.

But the formula indicates the governing character of the minimum crack growth rate for the load sequence effect. This minimum is a important as the length of the interaction zone, $\triangle a_D$. Thanks to the publication of min da/dN data in two references / 7, 8 / we are in a position to corroborate this finding before going forth. In Fig. 6 experimentally found numbers of overload affected cycles are depicted in a normalized manner versus the respective values of \uparrow obtained from the published minimum da/dN values. The normalization is made with the number

1)
$$da/dN = C \cdot \triangle K_{eff}^{m}$$

of C.A. cycles for the same increase in crack length. The materials are a mild steel and aluminum 2024 both with an exponent of m = 3, approximately. The associated theoretical curve runs in the midth of the experimental results. The outer lines are the factor of two lines with respect to the relative cycle number.

With this encouraging proof in mind we continue to seek for a closed solution where the minimum crack growth rate can be expressed through known parameter values. Fig. 7 demonstrates the way of thinking based on the findings of the Fatigue Fracture Mechanics studies /4. They revealed that the development of the plastic zone size with crack growth – see the lower right hand corner – is quite similar to the $\triangle K_{eff}$ development. The respective two minima just occur when the plastic zone strikes the compressive plastic zone of the peak cycle /3.

As, under SSY assumptions, ΔK_{eff} is found to be proportional to the square root of the plastic zone size ω , which holds even for the sequence effect considered here, the crack growth rate minimum, too, must be a power function of the minimum plastic zone size. The proportionality factor can be obtained by fulfilling the limit condition with no overload. Hence, by inserting the closed SSY solution for the ω -minimum

$$\omega_n = \frac{2}{3} \left(\frac{\Delta K_E}{2G_Y} \right)^2 \cdot \left(P^* - P \right)^2$$

$$\Delta K_E = K_E \cdot K_G$$

into y of the integral solution

$$V = \frac{(da/dN)_D}{(da/dN)_M} = \left(\frac{P^X - 1}{P^X - P}\right)^M$$

the final formula is obtained. The intermediate steps of the derivation can be made straight forwardly by using the solution, the authors presented at the ASTM Fracture Symposium /9.

The one point to feature out is, that the plastic zone size minimum is a function of what the authors call the "peak cycle ratio P", defined in the upper right hand corner of Fig. 7. This P-ratio turns out to be the independent variable in the relationship for the N_{SP} / N_{CA} - ratio.

The final formula is written down in the upper left hand corner of Fig. 8. P^* and τ have theoretically known values: $P^* = 2$, $\tau = m - 1$. Since, however, the real material behaviour could not be correctly modelled by the approach, these two quantities were taken as free constants. P^* is that value of P where crack arrest i.e. infinite crack growth life can be observed. Theoretically, it occurs if P equals P what, in case of P neans a hundred percent overload with no underload. However, experimental values lie between the values of P and P are a function of the P-ratio and the two material constants one straightforwardly gets the desired number of delay cycles for any load values considered.

Instead of only two references available for the first proof of the theoretical approach we can now use results of approximately five hundred experiments on six materials by twelve authors / 5, 10 - 20 /. Only few of them are shown in the figures.

A difficulty arose from the fact that mostly no constant amplitude data of the tested materials were published (cf. the suggestion in / 21 /). They had to be estimated. This fact must be taken into consideration when judging on the scatter of results. The symbols in Figures 9 and 10 represent the experimentally evaluated N_{SP}/N_{CA} ratio versus the peak range ratio of the load sequence applied. You also see the theoretical curve with the values which seem to fit the experimental data best. The comparison of the Aluminum to the Titanium data yields less delay for the Titanium. The exponents differ whereas the crack arrest ratio is identical for both materials.

On the left hand side of Fig. 10 the experimental results for three different steels are compared to the theoretical curve with the indicated parameter values. On the right hand side one test group was analyzed with respect to the occurence frequency. Only five of fiftyfive values lie outside the factor of 2 limits. These are mostly long life data with P values greater than 2 where the scatterband in cycles loses its usefulness.

To summarize the central idea of the paper:

The functional type of a relationship is presented describing single peak overload delay cycles as a function of a sole parameter. The absolute number of delay cycles depends on the C.A. crack growth behaviour of the material and at least two material constants. Those material constants must be evaluated from experiments. It is a first attempt to find the basic scheme behind this class of load sequences. Such a procedure would be promising also for other typical cases of load sequences. Valuable experiments are already carried out on this field but a similar interconnection between theory and experiment is not yet found. The present relationships are derived on the basis of elasto-plastic Fatigue Fracture Mechanics solutions including crack closure but with some very simplifying assumptions.

The authors wish to acknowledge the support given by the Deutsche Forschungsgemeinschaft (DFG) for important parts of this investigation.

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