

Deformation and Failure of Graphene Sheet and Graphene-Polymer Interface

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Abstract With a monolayer honeycomb-lattice of sp^2 -hybridized carbon atoms, graphene has demonstrated exceptional electrical, mechanical and thermal properties. One of its promising applications is to create graphene-polymer nanocomposites with tailored mechanical and physical properties. In general, the mechanical properties of graphene nanofiller as well as graphene-polymer interface govern the overall mechanical performance of graphene-polymer nanocomposites. However, the strengthening and toughening mechanisms in these novel nanocomposites have not been well understood. In this work, the deformation and failure of graphene sheet and graphene-polymer interface were investigated using molecular dynamics (MD) simulations. The effect of structural defects on the mechanical properties of graphene and graphene-polymer interface was investigated as well. The results showed that structural defects in graphene (e.g. Stone-Wales defect and multi-vacancy defect) can significantly deteriorate the fracture strength of graphene but may still make full utilization of corresponding strength of graphene and keep the interfacial strength and the overall mechanical performance of graphene-polymer nanocomposites.

Keywords Graphene, structural defect, fracture strength, graphene-polymer interface, load transfer

1. Introduction

Graphene has attracted increasing research effort since its discovery [1], largely due to its excellent electrical, mechanical and thermal properties. For example, graphene has high electron mobility ($25000 \text{ cm}^2/\text{Vs}$) at room temperature [1], anomalous quantum Hall effect [2], extremely high Young's modulus ($\sim 1 \text{ TPa}$) and fracture strength ($\sim 130 \text{ GPa}$) [3] and superior thermal conductivity ($5000 \text{ Wm}^{-1}\text{K}^{-1}$) [4]. These exceptional properties make graphene an ideal candidate as reinforcement in functional and structural polymer composites. For instance, graphene-polymer composite has a electrical conductivity of $\sim 0.1 \text{ Sm}^{-1}$ when adding only 1 v% graphene [5]. Poly(acrylonitrile) with 1 wt% functionalized graphene obtains a remarkable shift in glass transition temperature of over $40 \text{ }^\circ\text{C}$ [6]. More significantly, the composites show notable improvement in fracture strength and toughness, buckling and fatigue resistance [7-12].

Graphene can be produced via chemical vapour deposition (CVD) [13], mechanical exfoliation [14], chemical reduction of graphene oxide sheets [15], etc. It has been confirmed the properties of graphene can be modified by chemical functionalization [16-18]. However, both material production processes and chemical treatment may introduce structural defects in graphene, such as Stone-Wales (S-W) type defects (nonhexagonal rings generated by reconstruction of graphene lattice) [19], single and multiply vacancies, dislocation like defects, carbon adatoms, or accessory chemical groups. Recently, Gorjizadeh et al. [20] demonstrated that the conductance decreases in defective graphene sheets. Pei et al. [21, 22] studied the influence of functionalized groups on mechanical properties of graphene. Furthermore, it is still not well understood the underlying strengthening and toughening mechanisms of graphene-polymer nanocomposites and the influence of defective graphene on them. Further study is much required. Due to the nano-scale dimensions, it is difficult to accurately evaluate the properties of graphene sheets via experiment. Alternatively, molecular dynamics (MD) method has been widely utilized to investigate carbon-based nanomaterials [23-27]. In this work, we present a MD investigation on the fracture strength of graphene with structural defects (S-W defect and multi-vacancy defect) and interfacial behaviour of

graphene-polymer nanocomposites. Our results are helpful for a better understanding of the design and performance of graphene-polymer nanocomposites.

2. Models and Methods

To simulate a monolayer graphene sheet, a MD model ($42.6 \text{ \AA} \times 49.2 \text{ \AA}$) was built that consists of 800 carbon atoms. As confirmed by Zhao et al. [28], the possible model size effect on mechanical properties can be largely neglected when the diagonal length is over 5 nm. Therefore, the diagonal length of our model was chosen as 6.51 nm. The model was firstly relaxed to a minimum energy state with the conjugate gradient energy minimization. Then, Nose-Hoover thermostat [29, 30] was employed to equilibrate the graphene sheet at a certain temperature with periodic boundary conditions (PBCs). The adaptive intermolecular reactive bond order (AIREBO) potential [31] implemented in the software package LAMMPS [32], was used to simulate covalent bond formation and bond breaking. Such AIREBO potential has successfully simulated and predicted mechanical properties of carbon-based materials, i.e. fullerene, carbon nanotube and graphene.

In order to prepare the atomistic structures of graphene-polymer nanocomposite for simulation, a two-dimensional (2D) periodic model of polyethylene (PE) layer independent of graphene nanofiller was established. The polymer system consists of 25 PE molecules, with each molecule ($\text{CH}_3\text{-(CH}_2\text{-CH}_2\text{)}_{59}\text{-CH}_3$) composed of about 60 monomers. All the PE chains were prepared by commercial software Material Studio developed by Accelrys Inc. Then, two types of graphene-PE unit cells were constructed by stacking two PE layers with single graphene sheet (Case 1) and single defective graphene sheet (Case 2). To investigate the interfacial characteristics of graphene-PE nanocomposite, an ab initio force field polymer consistent force field (PCFF) [33, 34] was employed with the effective open-source code LAMMPS [32]. The interfacial interaction has been widely investigated in carbon-based materials and polymer-matrix nanocomposites [35-41]. To obtain the equilibrated structure of such unit cell, the model was first put into a constant-temperature, constant-pressure (NPT) ensemble for 250 ps by fixing the graphene with temperature of $T=100 \text{ K}$, pressure of $P=1 \text{ atm}$ and time step of $\Delta t=1 \text{ fs}$ after initial energy minimization (stage 1). Then, the unit cell model was further equilibrated for 250 ps with the same NPT ensemble and time step (stage 2). For the pull-out simulation of graphene from PE matrix, the displacement increment along the x axis of $\Delta x=0.001 \text{ \AA}$ was applied. During such process, graphene nanofiller were fixed while PE matrix was relaxed to equilibrate the whole dynamic system.

3. Results and Discussion

3.1 Effect of Stone-Wales (S-W) on Fracture Strength

Due to the short-ranged covalent bonding between carbon atoms, bond rotation and bond breaking are two basic deformation mechanisms in graphene. In this study, we considered two types of S-W defects, namely S-W_1 and S-W_2 , which are caused by 90° rotation of C-C bonds in different directions, as shown in Figure 1(a). Figure 1(b) showed corresponding stress-strain curves at different temperatures (300 K~900 K). In terms of true (Cauchy) stress, fracture strengths along armchair and zigzag directions at 300 K are 104 and 127 GPa, respectively. These values are in good agreement with experiment results $\sigma_f \sim 130 \text{ GPa}$ [3], as well as previous atomistic simulation results [26, 28].

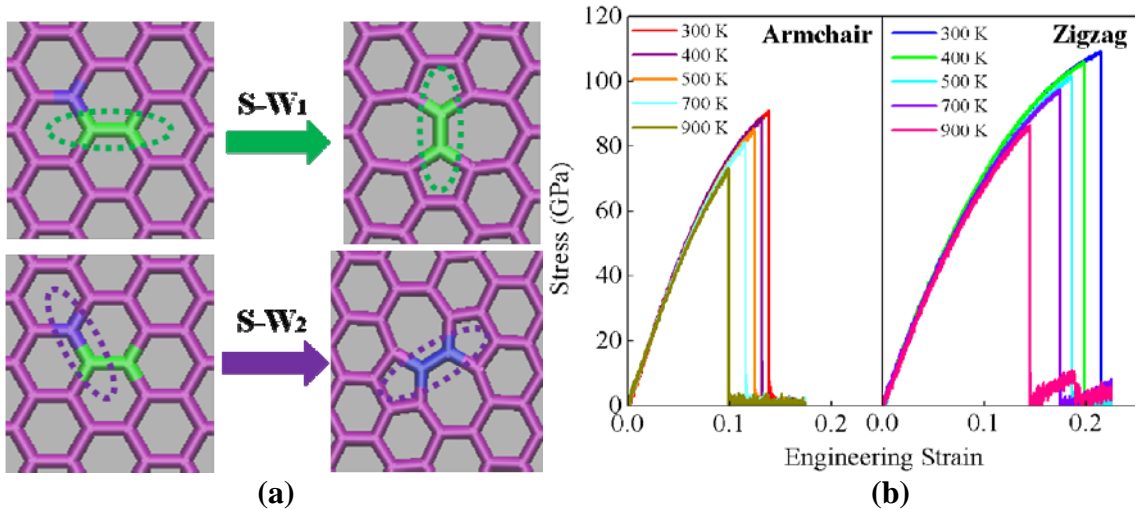


Figure 1. (a) Atomistic models of S-W₁ and S-W₂ defects. (b) Stress-strain curves of pristine graphene sheet under uniaxial tension along armchair and zigzag directions at different temperatures (300~900 K).

Our MD simulation results demonstrated that both S-W defects and temperature could significantly deteriorate the fracture strength of graphene. Figure 2 showed the fracture strength of graphene with S-W defects at different temperatures (300 K ~ 900 K). Under zigzag loading, the average strength loss caused by S-W₁ and S-W₂ was 16.26% and 41.30%, respectively. Under armchair loading, the average loss of fracture strength by S-W₂ defect was about 15.75%. For S-W₁, however, fracture strength increased when temperature was above 600 K (C point in Figure 2(a)). This was attributed to the healing of S-W₁ defect with increasing temperature. As shown in Figure 2(b), at 600 K, the S-W₁ defect was stable. At 700 K, however, the S-W₁ defect was healed by 90° rotation of C-C bond. As mentioned above, mechanical strain could lower the healing energy barrier E_{eb}^- . Therefore, according to the kinetic rate of the healing of S-W defects (ν),

$$\nu = af_0 \exp^{-E_{eb}^-/kT} \quad (1)$$

Where f_0 is attempt frequency (about $10^{13}/s$); k is the Boltzmann's constant and a is the lattice spacing $a = \sqrt{3}r_0$, where $r_0 = 1.42 \text{ \AA}$ is the C-C bond length. From Eq. 1, the healing of S-W₁ defect became easier with increase of mechanical strain and temperature, consistent with the MD simulation.

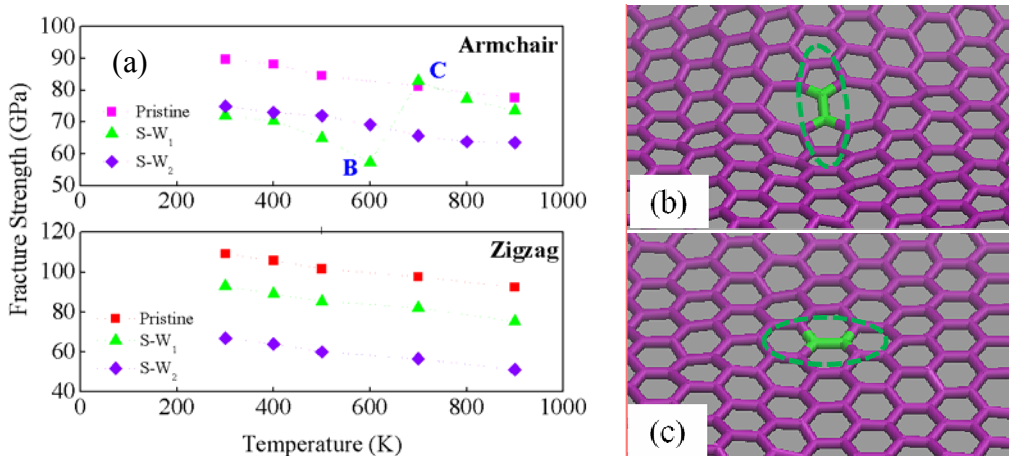


Figure 2. (a) Fracture strength of pristine, S-W₁ and S-W₂ defected graphene versus temperature under armchair and zigzag loading conditions. (b-c) Configuration change in the pre-existing S-W₁

defect from point B (b) (at 600 K) to point C (c) (at 700 K) highlighted in (a).

3.2 Effect of Vacancy Defect on Fracture Strength

In the temperature range of 300~900 K, the fracture strength of graphene sheet with vacancy was evaluated under tension along the armchair direction. The simulation model with 1, 2, and 3 vacancies is shown in Figure 3(a). Figure 3(b) showed the fracture strength σ_f for the graphene sheets with different vacancy number at temperatures 300 K, 500 K, and 900 K. It can be seen that fracture strength decreases with increasing temperature as well as the number of vacancy. For the sheet with 3 vacancies, the fracture strength loss was 37.3%, 40.2% and 42.4%, corresponding to 300, 500 and 900 K, respectively. Therefore, atomic scale defect such as vacancy does play a critical role in dictating the mechanical performance of graphene.

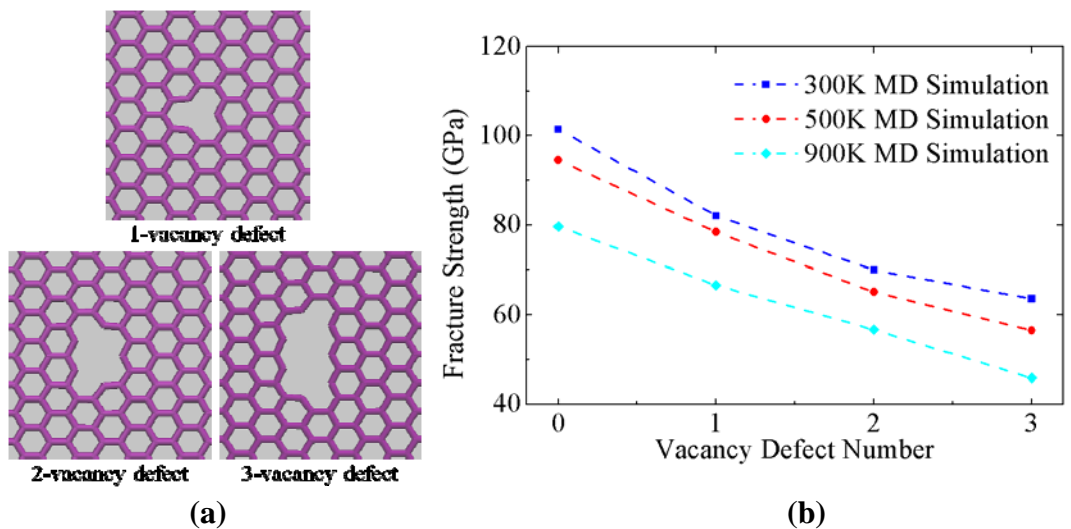


Figure 3. (a) Three types of vacancy defect. (b) Fracture strength of defective graphene sheet versus the number of vacancy defect.

3.3 Interfacial Behaviour of Graphene-PE Nanocomposites

To study the interfacial behavior of graphene-PE nanocomposites, MD simulation of pull-out test was carried out by pulling out the graphene nanofiller from PE matrix. In particular, interfacial shear force (ISF) can govern the effectiveness of load transfer during pull-out process. On the basis of the expressions $f_{ISF} = -\partial E_{int} / \partial X$, f_{ISF} can be calculated in terms of the given $E_{int}-X$ curve. As shown in Figure 4(a), $f_{ISF}-X$ curve can also be divided into three stages. It can be found that the Stage I and Stage III had approximately the same range of $X_I=X_{III}=1.0$ nm, which was close to the cut-off distance of vdW interaction. At Stage I, the magnitude of f_{ISF} rose quickly in all cases. The increase of f_{ISF} can be attributed to the newly formed surface of graphene (outside part). Then, f_{ISF} went through a long and approximate platform at Stage II. This was because that the length of the effective newly formed surface kept at 1 nm from the pull-out end, thus leading to constant f_{ISF} . Finally, it decreased to the value similar to that at first beginning until complete pull-out. At all three stages, the values of f_{ISF} varied periodically with the variation of X , largely due to the inhomogeneous distribution of PE structures in the interfacial area. The effect of structural defect on f_{ISF} was considered as well. However, the effect of structural defects on f_{ISF} is unobvious, which might be beneficial for application of graphene-polymer nanocomposites. Then, pull-out stress could be calculated in the expression as $\sigma_p = f_{ISF} / A_{eff}$, where A_{eff} is the effective cross section of graphene sheet, $A_{eff} = Wt$ with $W = 50$ Å the width of graphene sheet and $t = 3.44$ Å the sheet thickness.

According to Figure 4, the maximum pull-out stress σ_p^{\max} in both cases was about 0.87 GPa, much lower than fracture strength of pure graphene and defective graphene, namely $\sigma_p^{\max} / \sigma_f = 1$. Therefore, in the practical applications of graphene-polymer nanocomposites, graphene nanofiller could be full utilized even with structural defects on its surface. However, increasing the utilization of fracture strength of graphene deserves further investigation in future.

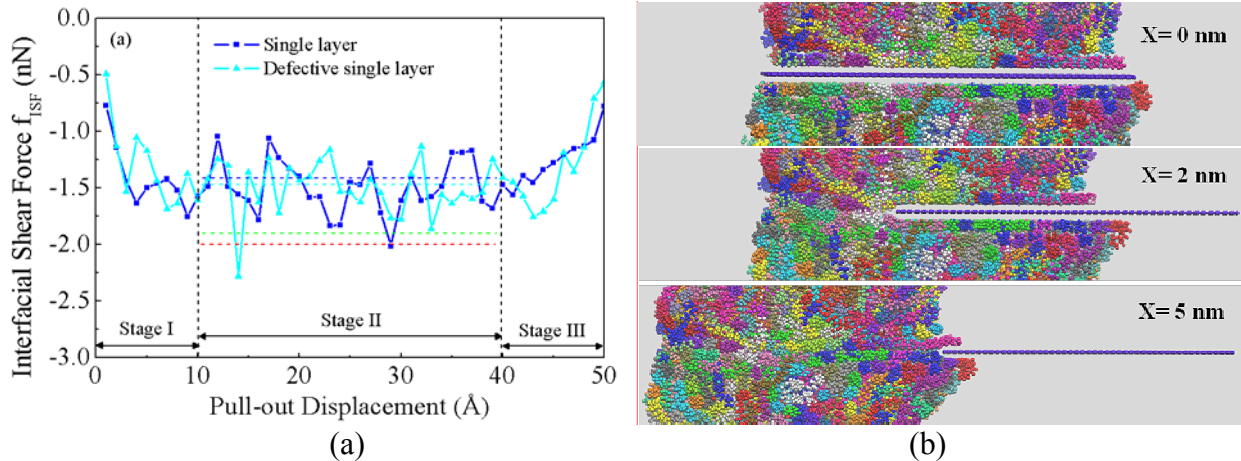


Figure 4. (a) Interfacial shear force f_{ISF} versus pull-out displacement X in Case 1 (pure graphene sheet) and Case 2 (defective graphene sheet). (b) Snap shots of graphene pull-out from PE matrix in Case 1.

4. Conclusions

Deformation and failure of graphene sheet and graphene-polymer interface was investigated using MD simulations. Calculation results showed that the fracture strength of graphene is dependent on both structural defects and temperature. Structural defects in graphene (e.g. Stone-Wales defect and multi-vacancy defect) can remarkably damage the fracture strength of graphene. As for the interfacial behaviour of graphene-polymer nanocomposites, it was shown that the value of ISF vary at each end of the graphene nanofiller within the range of 1 nm, while keep approximately constant (ISF) and zero (ISS) at middle stage. Particularly, ISF is independent of the pre-existing vacancy defect in graphene. Furthermore, graphene nanofiller could be full utilized even with structural defects owing to much lower pull-out stress. The study of the mechanical performance of defective graphene and graphene-polymer interface sheds light on the better understanding of design and application of graphene-polymer nanocomposites.

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