

Studies on a Three-dimensional Semi-infinite Crack

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ABSTRACT

The purpose of this paper is to report on a new approach to deal with the analytical evaluation for the solution of a three dimensional problem of a semi-infinite stress-free crack in an unbounded, linearly, elastic, isotropic medium.

The fundamental solution of a Green's function for a point load acting in the interior of an unbounded, homogeneous, isotropic, thick layer is used in connection with the boundary integral equations technique in order to obtain a formulation for the stress and displacement fields around the crack. This new formulation is based in the Green function obtained by the author and defined elsewhere. The formulation may analogously be applied to other three-dimensional cracked finite or infinite configurations.

KEYWORDS

Plate solution, 3-D Crack, Boundary Element Method.

1. REVIEW OF THE BOUNDARY INTEGRAL EQUATION METHOD IN THREE-DIMENSIONAL CRACKED CONFIGURATIONS

The boundary-integral equation method for elastic stress analysis is now well established as a complementary tool to finite element methods (Banerjee, 1981; Brebbia, 1978). Some of the advantages over other numerical solution procedures are well known. Foremost is the reduction of the dimensionality of the problem by one. Secondly, it is specially well suited to problems extending to infinity. Thirdly, since the formulation is based on fundamental solutions that satisfy the governing differential equations, approximation of the variables is required only on the boundary in the numerical solution of the equations. Once the boundary solution has been numerically obtained, interior values may easily be determined. These features are particularly advantageous for modeling regions with high stress gradients with great accuracy and efficiency, making this technique an appealing tool for numerical solution of problems in linear-elastic fracture mechanics.

In a number of papers, various boundary integral formulations have been shown to be useful for particular classes of boundary-value problems. The main difference among them focuses on the use of diverse Green's functions which are more appropriate for dealing with the geometry of the problem considered. Thus, for generic two-dimensional and three-dimensional elastostatic problems, Kelvin's fundamental solution of a concentrated

load in an infinite medium has been widely used (Risso, 1967; Cruse, 1969). Also, for problems involving a free-surface the solution presented by Melan (1932), for the stress distribution due to point loads applied within the isotropic half-space, or the one given by Mindlin (1936), for the half-space are of an utter interest. These fundamental solutions have been applied to the boundary element technique (BEM) by Telles & Brebbia (1981).

Further, for elastostatic problems containing cutouts, Kelvin's solutions in 2D and 3D have been profusely used. On this lane rests the pioneering work of Cruse and VanBuren (1971). Since then a huge stack of papers involving cracks, tackled by the boundary integral technique, have been published. Among them, the early study by Snyder and Cruse (1973) is worth to be mentioned. In their report several finite anisotropic two-dimensional plates were analysed by using the Green function corresponding to a concentrated load applied to an infinite plane containing a semi-infinite stress free crack. Thus, two-dimensional crack problems with finite geometry are then solved by the BEM with no crack modeling required; the crack presence being accounted for by the new fundamental solutions.

Several numerical formulations have been devised for the application of the boundary integral equation method to fracture mechanics problems in three-dimensions. Pioneering work, in this field, was that of Cruse (1970, 1972, 1973, 1974, 1975, 1977, 1971, 1977) and Tan & Fenner (1978), among others (1985).

It is the purpose of this paper to report on the application of a new analytical method to the solution of the three-dimensional through crack on an infinite plate, as well as its use to other 3-D cracked, finite or infinite configurations.

2. INTEGRAL REPRESENTATIONS

This section deals with the integral formulation of the elastostatic problem for the layer-space defined for the case of a semi-infinite three-dimensional crack. The following developments are based on previous progress, by the author (1985, 1987, 1987), in the development of the necessary theoretical background for the determination of stresses and displacements in three-dimensional plates subjected to concentrated point loads.

Consider a linear, elastic, isotropic, homogeneous medium occupying a space described by an infinite three-dimensional thick layer subjected to a concentrated body force acting at a point ξ in a direction indicated by the unit vector e^j and having the form

$$F_i(\mathbf{x}) = \delta(\mathbf{x} - \xi) \delta_{ij} e^j,$$

where \mathbf{x} is a point of the body, $\delta(\cdot)$ is the Dirac delta function and δ_{ij} is Kronecker's delta. The response of this body, obtained as the solution of the equations of elastostatics is given by the displacement vector and the stress tensor

$$\begin{aligned} u_i(\mathbf{x}) &= u_i^j(\mathbf{x}; \xi) e^j, \\ \sigma_{ik}(\mathbf{x}) &= \sigma_{ik}^j(\mathbf{x}; \xi) e^j, \end{aligned}$$

where the second-order symmetric tensor $u_i^j(\mathbf{x})$ and the third-order tensor σ_{ik}^j (fundamental singular solution) express the displacement component in the i -direction and the stress ij -component, respectively, at the point \mathbf{x} due to a concentrated force of magnitude F acting at the point ξ in the j -direction.

The fundamental solution pair $[u^j(\mathbf{x}; \xi), \sigma^j(\mathbf{x}; \xi)]$, defined for all points \mathbf{x} except ξ , is called the three-dimensional Layer-state and is characterized by the properties established elsewhere.

The boundary integral equation technique is based on the Betti reciprocal work theorem, which provides an integral relationship between two elastostatic states. If one of them corresponds to a fundamental singular solution, Somigliana's identity is obtained. Depending on the fundamental solution used, various formulations may be available and more suitable for solving certain problems.

Consider now a linear, elastic, isotropic, homogeneous body R , with surface ∂R , subjected to body forces $F(\mathbf{x})$, surface tractions $s(\mathbf{x})$ and boundary displacements $u(\mathbf{x})$ (boundary conditions of Neumann, Dirichlet or mixed type). Using the three-dimensional Layer-state, one can derive Somigliana's integral representation for the displacement and stress field, in the form

$$\begin{aligned} u_j(\xi) &= - \int_{\partial R} s_i^j(\mathbf{x}; \xi) u_i(\mathbf{x}) dA_s + \int_{\partial R} u_i^j(\mathbf{x}; \xi) s_i(\mathbf{x}) dA_s \\ &\quad + \int_R u_i^j(\mathbf{x}; \xi) F_i(\mathbf{x}) dV_s, \end{aligned} \quad (1)$$

$$\begin{aligned} \sigma_{ij}(\xi) &= - \int_{\partial R} \bar{s}_k^{ij}(\mathbf{x}; \xi) u_k(\mathbf{x}) dA_s + \int_{\partial R} \bar{u}_k^{ij}(\mathbf{x}; \xi) s_k(\mathbf{x}) dA_s \\ &\quad + \int_R \bar{u}_k^{ij}(\mathbf{x}; \xi) F_k(\mathbf{x}) dV_s, \end{aligned} \quad (2)$$

where ξ is a point of the body, \mathbf{x} a point in the body or on the surface; u , s , F are the displacement, traction and body force vectors, respectively; u^j , σ^j , u^{ij} , σ^{ij} are the Green tensors as described elsewhere.

The solution of the elastostatic problem can be inferred from the boundary equation, derived from (1)

$$\begin{aligned} c_{ji}(\xi) u_i(\xi) &+ \int_{\partial R} s_i^j(\mathbf{x}; \xi) u_i(\mathbf{x}) dA_s \\ &= \int_{\partial R} u_i^j(\mathbf{x}; \xi) s_i(\mathbf{x}) dA_s + \int_R u_i^j(\mathbf{x}; \xi) F_i(\mathbf{x}) dV_s, \end{aligned} \quad (3)$$

where the tensor $c_{ji}(\xi)$ is equal to 1/2 if ξ is at a smooth surface.

In problems with zero body forces, the last integral in equations (1), (2) and (3) vanish. Thus, equation (3) becomes

$$c_{ji}(\xi) u_i(\xi) + \int_{\partial R} s_i^j(\mathbf{x}; \xi) u_i(\mathbf{x}) dA_s = \int_{\partial R} u_i^j(\mathbf{x}; \xi) s_i(\mathbf{x}) dA_s, \quad (4)$$

which relates the boundary displacements u and tractions s .

For arbitrary geometries and complicated variations of the related functions, equation (3) or (4) has to be treated numerically.

3. BOUNDARY EQUATION FOR A SEMI-INFINITE THREE-DIMENSIONAL THROUGH-CRACK

Consider a three-dimensional infinite layer R containing a traction free crack L . Assume that the region R is embedded in an infinite plate and that a unit load $F(\theta)$ is applied at some point θ in the region R but not on the boundary ∂R . The system under consideration and the coordinate axes are depicted in Fig.1.

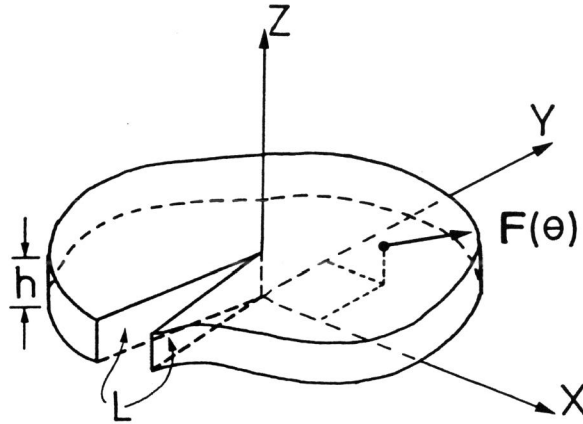


Fig. 1. Three-dimensional infinite layer containing a semi-infinite through-crack subjected to a unit load applied at point θ .

If the load point is represented by a unit force in the k -direction,

$$F_i(\mathbf{x}) = \delta(\mathbf{x} - \theta) \delta_{ik} e^k, \quad (5)$$

the displacement and traction field will be given by

$$u_i(\mathbf{x}) = \hat{u}_i^k(\mathbf{x}; \theta) e^k, \quad (6)$$

$$s_i(\mathbf{x}) = \hat{s}_i^k(\mathbf{x}; \theta) e^k, \quad (7)$$

where $\hat{u}^k(\mathbf{x}; \theta)$ and $\hat{s}^k(\mathbf{x}; \theta)$ represent the corresponding solution vectors for displacement and traction, respectively, for a concentrated load parallel to the k -axis.

Let S_u , S_l denote the upper and lower planar surfaces of the plate, respectively, and L the crack surface; thus $\partial R = S_u \cup S_l \cup L$. Since the two planar surfaces are traction free, then

$$\int_{S_u \cup S_l} u_i^k(\mathbf{x}; \xi) s_i(\mathbf{x}) dA_s = 0,$$

and the second integral in (3) is exclusively extended to the surface crack.

Also, the first integral in (3) provides

$$\int_{S_u \cup S_l} s_i^k(\mathbf{x}; \xi) u_i(\mathbf{x}) dA_s = 0,$$

as s^j vanishes there (traction free condition of the analytical point load solution).

The third integral in (3), taking into account (5), has the form

$$\int_R u_i^k(\mathbf{x}; \xi) \delta(\mathbf{x} - \theta) \delta_{ik} e^k dV_s = u_i^k(\theta; \xi) e^k.$$

From the above and (4.2), expression (3) results in

$$c_{ji}(\xi) \hat{u}_i^k(\xi; \theta) = - \int_L s_i^k(\mathbf{x}; \xi) \hat{u}_i^k(\mathbf{x}; \theta) dA_s + u_i^k(\theta; \xi),$$

where summation is implied on the index i .

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